

10/9 Lecture outline

- Write  $U = U(T, V)$ . Exact differential means

$$dU = \left(\frac{\partial U}{\partial T}\right)_V dT + \left(\frac{\partial U}{\partial V}\right)_T dV, \quad \text{with} \quad \left(\frac{\partial}{\partial V} \left(\frac{\partial U}{\partial T}\right)_V\right)_T = \left(\frac{\partial}{\partial T} \left(\frac{\partial U}{\partial V}\right)_T\right)_V.$$

Using first law,

$$C_V = \left(\frac{\partial U}{\partial T}\right)_V, \quad C_P = C_V + \left[\left(\frac{\partial U}{\partial V}\right)_T + P\right] V\beta.$$

Note that  $c_P$ ,  $c_V$ , and  $\beta$  are easily measured, as are  $P$  and  $V$ . So we can use the above to then determine  $U(T, V)$ .

- Physically clear that  $C_P > C_V$ : consider fixed  $dQ$  heating two containers, one with  $V$  fixed and one with  $P$  fixed. Get  $(dW)_V = 0$  and  $(dW)_P > 0$ . So  $(dU)_V > (dU)_P$ . Since  $T$  is a measure of  $U$ , get  $(dT)_V > (dT)_P$ ; so  $C_P > C_V$ . Let's show it mathematically. Use above to write

$$C_P - C_V = \left(\frac{dQ}{dV}\right)_T \left(\frac{\partial V}{\partial T}\right)_P = \frac{VT\beta^2}{\kappa_T}.$$

The last equality uses a Maxwell relation,  $\left(\frac{\partial P}{\partial T}\right)_V = T^{-1} \left(\frac{dQ}{dV}\right)_T$ , which is related to the fact that we'll see later that  $dQ = TdS$  (explain it a bit, via  $F^T = U - TS$ ). There is a stability condition, that systems find the state of lowest energy, which implies that  $\kappa_T > 0$ . The above relation then implies that  $C_P > C_V$ , as we expected. Define  $\gamma \equiv C_P/C_V$ .

- Plot of  $C_V$  and  $C_P$  over big temperature range. Interesting behavior at low  $T$ . Rather universal behavior at large  $T$  (ideal gas).

- Ideal gas has  $U = U(T)$ , so  $dQ = C_V dT + pdV$ . Using ideal gas law, get  $dQ = (C_V + nR)dT - VdP$ . Conclude  $C_P = C_V + nR$  for ideal gas. Thus  $C_V = nR/(\gamma - 1)$ ,  $C_P = \gamma nR/(\gamma - 1)$ .