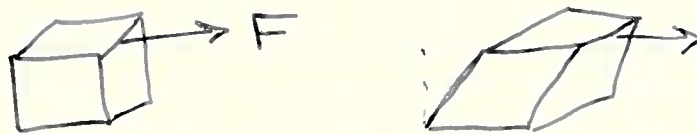


Fluids

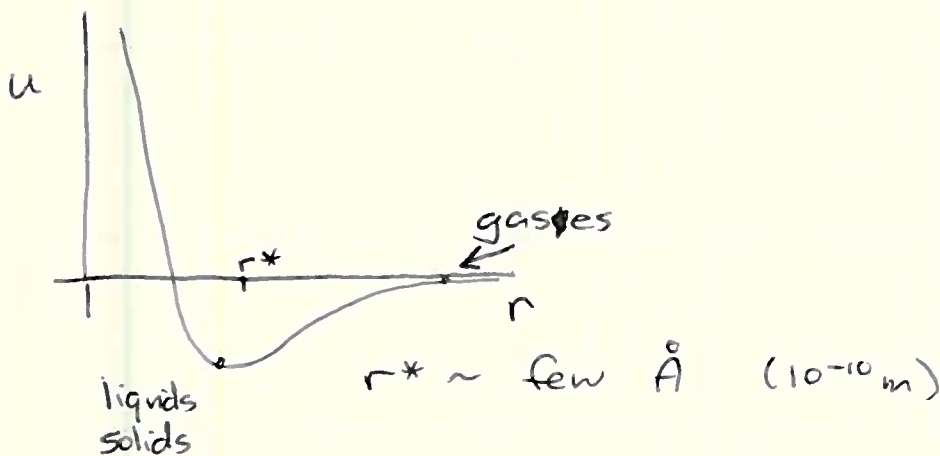
(vs solids) they flow in response to shear stresses



No restoring force.

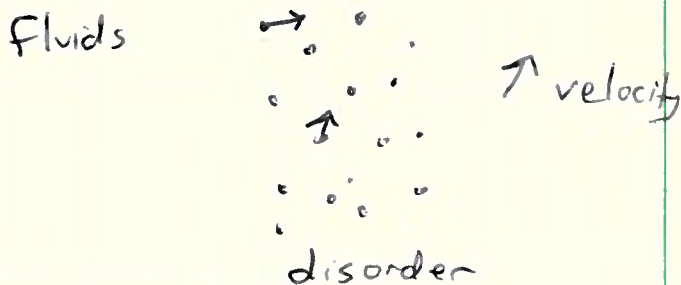
Do have restoring force for forces \perp to face.

On micro level, inter-molecule potential



lattice - order

shear restoring force \rightarrow preserves order



no memory of position

Density

$$\rho = \frac{dm}{dv} \xrightarrow{\text{uniform}} \frac{m}{V} \quad [\rho] = \text{kg/m}^3$$

e.g.	best lab vac		10^{-17}	kg/m ³	
air	20°C + 1atm	1.21			\rightarrow compressible
	20°C + 50atm	60.5			
water	20°C + 1atm	$.998 \times 10^3$			\rightarrow incompressible
	20°C + 50atm	1.000×10^3			

black hole $\sim 10^{19}$

13,782 500 SHEETS FILLER 5 SQUARE
 43,381 50 SHEETS CYE-EASER 5 SQUARE
 43,382 100 SHEETS CYE-EASER 5 SQUARE
 43,383 200 SHEETS CYE-EASER 5 SQUARE
 43,384 100 SHEETS RECYCLED WHITE 5 SQUARE
 43,385 200 SHEETS RECYCLED WHITE 5 SQUARE
 Made in U.S.A.

Pressure

$$p = \frac{dF}{dA} \xrightarrow{\text{Uniform}} F/A$$

units $[p] = \text{Newton/meter}^2 \quad (\sim M/LT^2)$

$$1 \text{ Newton/meter}^2 \equiv 1 \text{ Pa} \quad \text{"pascal"}$$

e.g. $p_{\text{atmosphere}}$ (at sea level) $1 \times 10^5 \text{ Pa}$

$$\equiv 1 \text{ atm} = 14.7 \text{ lb/in}^2$$

Total force of atmosphere on palm of your hand:

$$A \approx 10 \text{ in}^2 \quad F = (14.7)(10) \approx 147 \text{ lbs}$$

Auto tire pressure (in excess of atmosphere pressure)

$$\text{Area } A \text{ of tire on ground} \sim (15 \text{ cm})^2 \times 4 \\ \sim .1 \text{ m}^2 \sim 160 \text{ in}^2$$

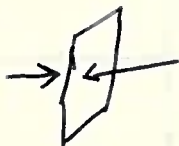
$$F = \text{weight of car} \sim 2 \times 10^3 \text{ kg} \times (9.8) \sim 2 \times 10^4 \text{ N} \sim 4000 \text{ lbs}$$

$$\therefore \text{pressure } p = 2 \times 10^5 \text{ Pa} \sim 25 \text{ lbs/in}^2$$

Hydrostatics: Fluids at rest. No shear forces, all forces \perp to every surface. Force element dF_{\perp} ,

on any surface element dA is

$$dF_{\perp} = -p \hat{n} dA \quad (\hat{n}: \text{outward normal vector})$$



(magnitude of F indep. of orientation)

$$\underline{p = \text{scalar}}$$

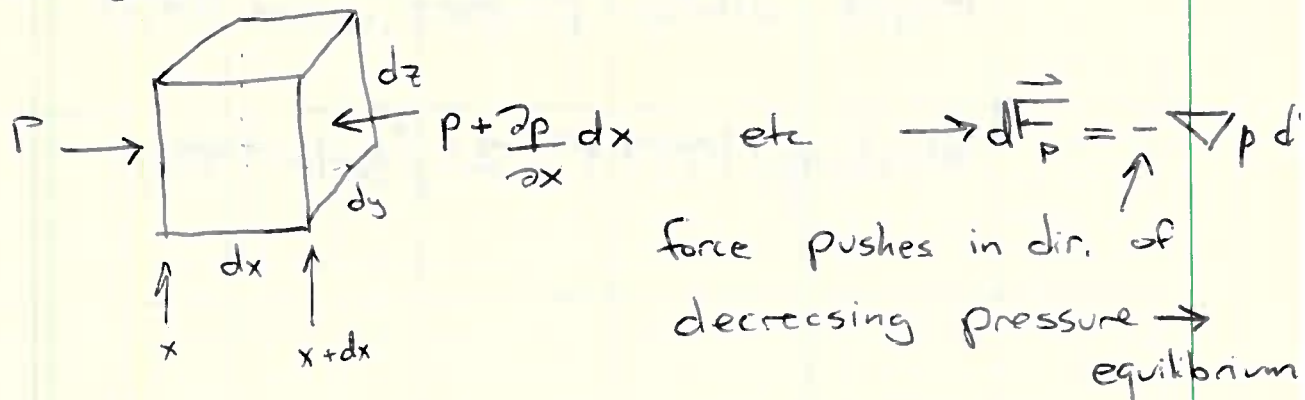
Pressure sensor:



$$F = k \Delta x = pA$$

$$p = k \Delta x / A$$

Force on small volume due to pressure



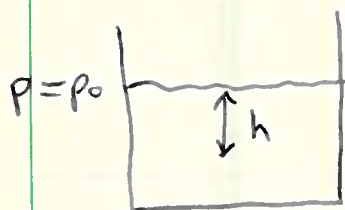
In equilibrium, with no outside forces, $p = \text{const}$

With gravity, $\vec{F}_{\text{total}} = \vec{F}_p + \vec{F}_{\text{grav}} = 0$

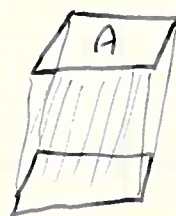
$$\vec{\nabla} p = \rho \vec{g} \quad \Rightarrow \quad p(z) = p(z_0) - \rho g (z - z_0)$$

(if $\rho \hat{=} \vec{g}$ const.)

\uparrow "incompressible"



$$p = p_0 + \rho g h$$



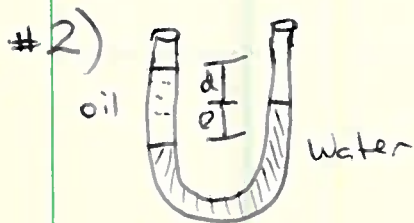
weight $\rho_0 V$

$$p_0 A + \rho h A c$$

Eg: 1) some waterproof watches quote max depth in atm
What is that in meters?

$$1 \text{ atm} = \Delta p = \rho_{\text{H}_2\text{O}} g \Delta h \quad \rightarrow \quad \Delta h = \Delta p / \rho_{\text{H}_2\text{O}} g$$

$$\Delta h = 1.01 \times 10^5 \text{ Pa} / (10^3 \text{ kg/m}^3) (9.8 \text{ m/s}^2) = 10.3 \text{ m} \sim 34 \text{ ft}$$

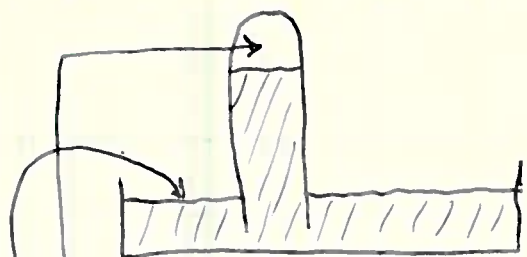


$$g(\rho_{\text{oil}}(d+l)) = g\rho_{\text{H}_2\text{O}} l$$

$$\Rightarrow \frac{d}{l} = \frac{\rho_{\text{H}_2\text{O}} - \rho_{\text{oil}}}{\rho_{\text{oil}}} \quad \text{indep. of } g$$

Measuring pressure with a Barometer:

- i) fill long glass tube with liquid
- ii) invert in a dish of the liquid



height h above level in dish

pressure here is $p \approx 0$ (almost vacuum)

pressure p here is $p_0 = \text{atmosphere pressure}$

equilibrium: $p_0 = \rho g h$

height h of barometer: $h = p_0 / \rho g$

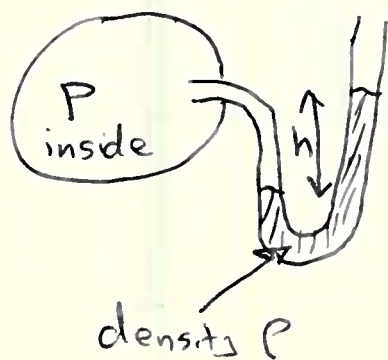
If liquid is water, using $p_0 = 1 \times 10^5 \text{ Pa}$

and $\rho_{\text{H}_2\text{O}} = 10^3 \text{ kg/m}^3$ $h \approx 10 \text{ meters}$ - too big!

Better to use mercury Hg, $\rho_{\text{Hg}} = 1.4 \times 10^4 \text{ kg/m}^3$

so now $h \approx .7 \text{ meters}$ - more manageable.

Open tube manometer: measure pressure difference
between fluid and atmosphere



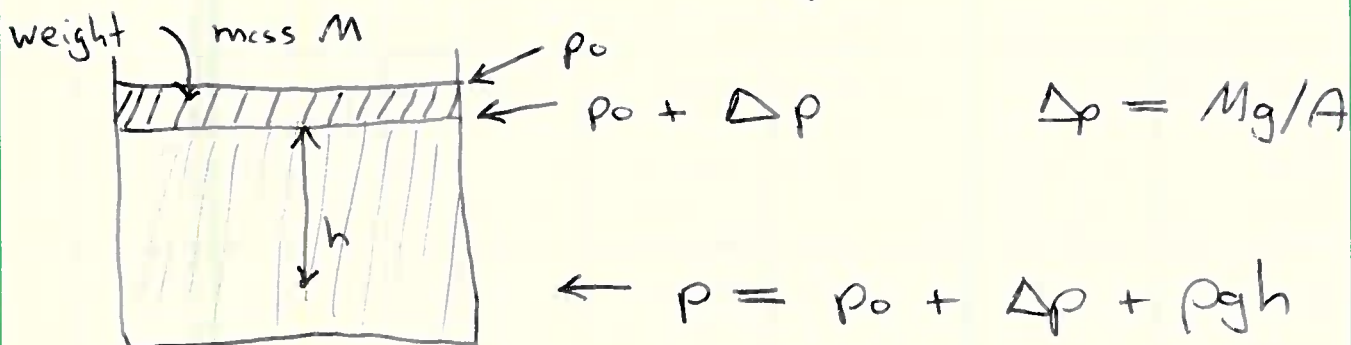
p_0 at open end

$$P = p_0 + \rho g h$$

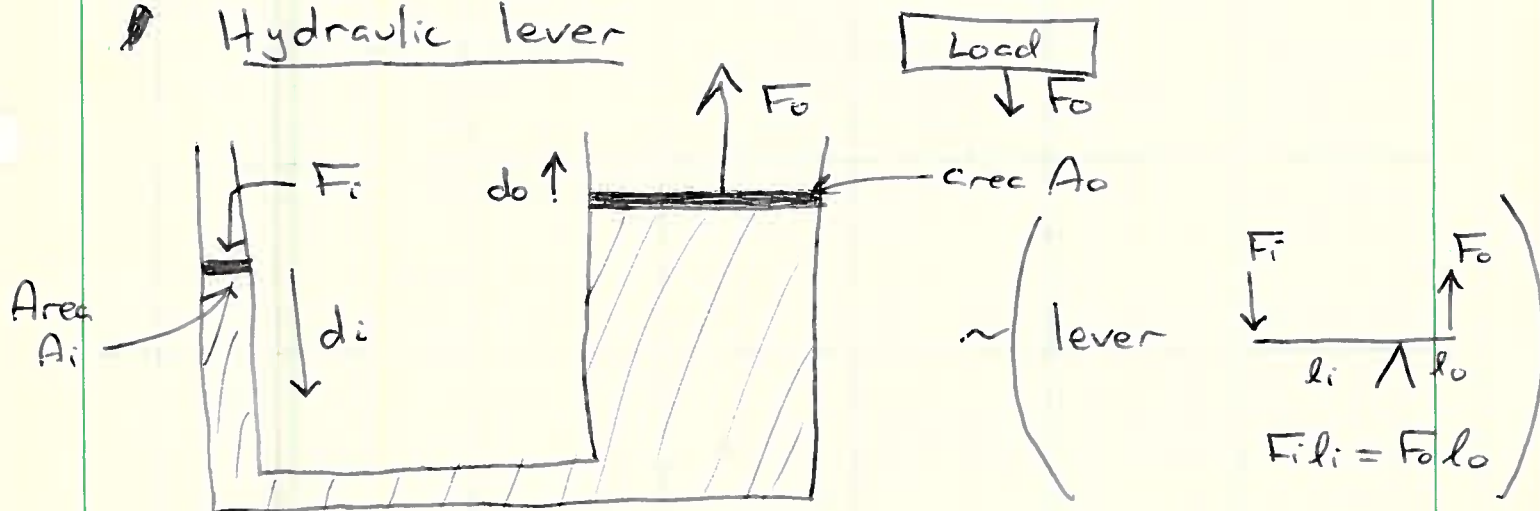
e.g. reading tire pressure.

Pascal's principle :

Additional pressure Δp_{ext} exerted on enclosed fluid (incompressible) $\rightarrow \Delta p_{\text{ext}}$ added to pressure everywhere in fluid.



Hydraulic lever



$$\Delta p = \frac{F_i}{A_i} = \frac{F_o}{A_o} \quad F_o = F_i \frac{A_o}{A_i}$$

magnified for $A_o > A_i$


Incompressibility : displacement $d_i A_i = d_o A_o$ (volumes same)

$$d_o = d_i A_i / A_o$$

Work done on left $F_i d_i = F_o d_o =$ work done by right.

no energy lost (to compression).

Archimedes' Principle: Pressure exerts force $d\vec{F} = -p \hat{n} dA$ on any area element. \hat{n} = outward normal

Consider object X inside of fluid 

Total force due to pressure = "buoyancy force"

$$\vec{F}_b = - \int p \hat{n} dA = - \int \vec{\nabla} p dV$$

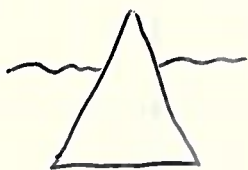
$\xrightarrow{\text{Surface of X}} S_x = \partial V_x$
 $\xrightarrow{\text{Volume of X}} V_x$

But in equilibrium $\vec{\nabla} p = -\rho_{\text{liquid}} g \hat{z}$

$$\text{so } \vec{F}_b = \hat{z} g \int_{V_x} dV \rho_{\text{liquid}} = \hat{z} \cdot (\text{Weight of displaced fluid})$$

$$\vec{F}_{\text{total}} = \vec{F}_b - M_x g \hat{z} = (M_{\text{displaced fluid}} - M_{\text{object}}) g \hat{z}$$

Eg Iceberg



$$\rho_{\text{ice}} = 917 \text{ kg/m}^3 @ 0^\circ\text{C}$$

$$\rho_{\text{sea water}} = 1024 \text{ kg/m}^3 @ 0^\circ\text{C}$$

$$F_b = g V_{\text{sub}} \rho_{\text{water}}, \quad V_{\text{sub}} = \text{submerged volume}$$

$$\text{Weight of iceberg: } W = g V_{\text{total}} \rho_{\text{ice}} = F_b$$

$$\Rightarrow V_{\text{sub}} \rho_{\text{water}} = V_{\text{total}} \rho_{\text{ice}}. \quad \text{Fraction above water:}$$

$$(V_{\text{total}} - V_{\text{sub}}) / V_{\text{total}} = 1 - \frac{\rho_{\text{ice}}}{\rho_{\text{H}_2\text{O}}} \approx 10\%$$

Hydrodynamics

$$\rho(\vec{r}, t), \quad p(\vec{r}, t), \quad \vec{v}(\vec{r}, t)$$

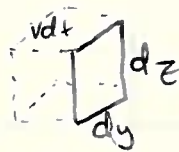
eqn of continuity:  mass in element time t : $\rho(\vec{r}, t) dV$

$$\Delta(\text{mass}) \text{ over time } dt : \frac{\partial \rho}{\partial t} dV dt = \Delta m$$

= mass which flows into region (conservation)

$$= -\nabla \cdot (\rho \vec{v}) dV dt$$

e.g. along x axis



amount $\rho v_x(x) dy dz dt$
flows in at x

$$\frac{1}{2} \text{ amount } \left[(\rho v_x)(x+dx) = \rho v_x(x) + \frac{\partial (\rho v_x)}{\partial x} dx \right] dy dz dt$$

flows out at $x+dx$.

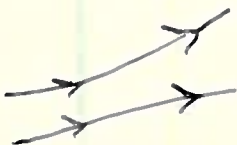
$$\text{so } \boxed{\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \vec{v}) = 0}$$

Here take $\rho = \text{const.}$: incompressible
above $\Rightarrow \nabla \cdot \vec{v} = 0$.

Steady flow : $\vec{v} = \vec{v}(\vec{r})$ no t dep.

non viscous : no friction or drag forces.

Stream lines : tangent to fluid velocity = actual path of fluid particles.

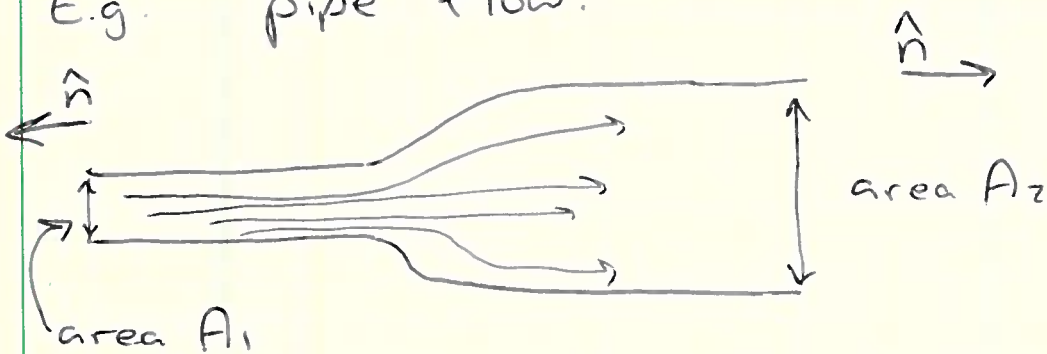


Integrate $\vec{\nabla} \cdot \vec{v} = 0$ over a volume V

$$0 = \int_V dV \vec{\nabla} \cdot \vec{v} = \int_{\partial V} \hat{n} \cdot \vec{v} dA$$

\hat{n} : normal to surface element dA

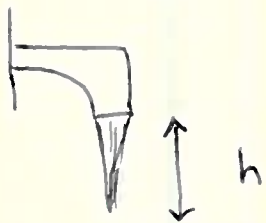
E.g. pipe flow:



$$0 = -v_1 A_1 + v_2 A_2 \quad \Rightarrow \quad v_1 A_1 = v_2 A_2$$

(larger area \rightarrow smaller velocity)

Faucet



$$v^2(h) = v_0^2 + 2gh$$

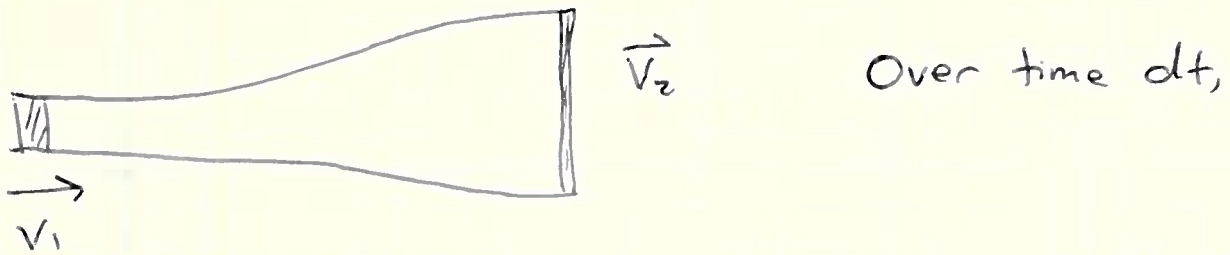
$$\pi a^2(h) v(h) = \pi a^2(0) v_0$$



$$\Rightarrow a(h) = a_0 / \left(1 + \frac{2gh}{v_0}\right)^{1/4}$$

Bernoulli's Equation: Energy conservation

Consider flow along streamline flow tube



Work done on fluid inside at end 1 = $p_1 A_1 v_1 dt$

Work done by fluid at end 2 = $p_2 A_2 v_2 dt$.

Net work done on fluid inside:

$$dW = p_1 A_1 v_1 dt - p_2 A_2 v_2 dt$$

continuity $\Rightarrow A_1 v_1 = A_2 v_2$ so $dW = (p_1 - p_2) A_1 v_1 dt$

Kinetic energy changes by $dK = \frac{1}{2} (\rho A_2 v_2 dt) v_2^2 -$

$$\frac{1}{2} (\rho A_1 v_1 dt) v_1^2 = \frac{1}{2} \rho (v_2^2 - v_1^2) A_1 v_1 dt.$$

Potential energy changes by: $du = (\rho A_2 v_2 dt) g h_2 -$

$$(\rho A_1 v_1 dt) g h_1 = \rho g (h_2 - h_1) A_1 v_1 dt.$$

Energy conservation: $dW = dK + dU \Rightarrow$

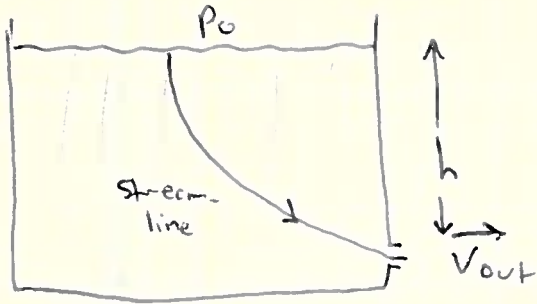
$$p_1 + \rho g h_1 + \frac{1}{2} \rho v_1^2 = p_2 + \rho g h_2 + \frac{1}{2} \rho v_2^2$$

ie.
$$p + \rho g h + \frac{1}{2} \rho v^2 = \text{constant.}$$

For $h = \text{const.}$, larger velocity \rightarrow lower pressure

e.g. blow between 2 pieces of paper.
They come together.

Eg: 1) Large tank with small hole



Large $\rightarrow V_{top} \approx 0$

$$(\text{@top}): P_0 + \rho gh = P_0 + \frac{1}{2} \rho V_{out}^2 \quad (\text{@hole})$$

$$\Rightarrow V_{out} = \sqrt{2gh}$$

same as for object dropped from height h.

more generally: velocity V_{top} with $V_t A_t = V_o A_o$

$$\rho gh + \frac{1}{2} \rho V_t^2 = \frac{1}{2} \rho V_o^2$$

$$V_o^2 \left(1 - \left(\frac{A_o}{A_t} \right)^2 \right) = 2gh \Rightarrow V_o = \frac{\sqrt{2gh}}{\sqrt{1 - \frac{A_o^2}{A_t^2}}}$$