6/5/17 Lecture 17 outline / summary

• Recall: $\mathcal{L} \supset \bar{\psi}(i\not{\!\!D} - m)\psi$, with $D_{\mu} = \partial_{\mu} + iqA_{\mu} + igT^{a}A_{\mu}^{a}$. $F_{\mu\nu} = [D_{\mu}, D_{\nu}]/(-ig) = \partial_{\mu}A_{\nu} - \partial_{\nu}A_{\mu} - ig[A_{\mu}, A_{\nu}]$, in the adjoint representation of the gauge group. $\mathcal{L} \supset -\frac{1}{4}TrF_{\mu\nu}F^{\mu\nu} \supset -gf^{abc}\partial_{\mu}A_{\nu}^{a}F^{\mu b}A^{\nu c} - (g^{2}/4)f^{abc}f^{ade}A_{\mu}^{b}A_{\nu}^{c}A^{\mu d}A^{\nu e}$.

• Continue with QCD Feynman rules, examples of color factors.

• Running couplings in QED vs QCD. Asymptotic freedom. Confinement. Motivate grand unification.

• On to the weak force! Two differences from other forces: it is chiral (hence parity violating), and the force carriers (Ws and Zs) are massive, which is why it is weak. "How can a force carrier be massive?" given that forces are related to gauge symmetries, and gauge invariance forbids mass terms (e.g. for the photon). Answer: the gauge invariance is *spontaneously broken* by the Higgs field. This is roughly similar to the Bose condensate in a superconductor.