

1/22/07 Lecture 4 outline

- Last time:

$$Z[J] = N \int [d\phi] e^{\frac{i}{\hbar}(S + \hbar \int d^4x J\phi)} = N \exp\left[\frac{i}{\hbar} S_{int}\left[-i \frac{\delta}{\delta J}\right]\right] Z_{free}[J],$$

where N is an irrelevant normalization factor (independent of J). Correspondingly, the green's functions are given by

$$\begin{aligned} G^{(n)}(x_1 \dots x_n) &= \frac{\int [d\phi] \phi(x_1) \dots \phi(x_n) \exp\left(\frac{i}{\hbar} S_I[\phi]\right) \exp\left[\frac{i}{\hbar} S_{free}\right]}{\int [d\phi] \exp\left(\frac{i}{\hbar} S_I[\phi]\right) \exp\left[\frac{i}{\hbar} S_{free}\right]}, \\ &= \frac{1}{Z[J]} \prod_{j=1}^n \left(-i\hbar \frac{\delta}{\delta J(x_j)}\right) \cdot Z[J]_{J=0}. \end{aligned}$$

(The denominator (in both lines) cancels off the vacuum bubble diagrams, which don't depend specifically on the Green's function.)

- Illustrate the above formulae, and relation to Feynman diagrams, e.g. $G^{(1)}$, $G^{(2)}$ and $G^{(4)}$ in $\lambda\phi^4$ theory. The functional integral accounts for all the Feynman diagrammer; even symmetry factors etc. come out simply from the derivatives w.r.t. the sources, and the expanding the exponentials.

- Recall story of cancellation of bubble diagrams. Recall for computing S-matrix elements, we will especially be interested in *connected* Green's functions. There are nice combinatoric formulae (you might have already seen some last quarter?). E.g.

$$\sum \text{all diagrams} = \left(\sum \text{“connected”}\right) \cdot \exp\left(\sum \text{disconnected vacuum bubbles}\right).$$

And the vacuum bubble diagrams cancel. We write “connected” because for $n > 2$ point functions there are still disconnected diagrams, connected to the external points, included in this sum.