

2/26/10 Lecture 16 outline

- New topic, functional integral for fermions and gauge fields. Path integral of same, general form. Need to understand some new issues with integrations. Fermions first. Grassmann number integrals, $\int d\theta(A + B\theta) = B$. Complex θ, θ^* , $\int d\theta^* d\theta \exp(-\theta^* b\theta) = b$. $\prod_i \int d\theta_i^* d\theta_i \exp(-\theta_i^* B_{ij} \theta_j) = \det B$. $\prod_i \int d\theta_i^* d\theta_i \exp(-\theta_i^* B_{ij} \theta_j) \theta_k \theta_l^* = (B^{-1})_{kl} \det B$.

- We can introduce sources for the fields:

$$\begin{aligned} Z[\bar{\eta}_i, \eta_i] &= \int d\bar{\theta}_i d\theta_i \exp(i(A_{ij} \bar{\theta}_i \theta_j + \bar{\eta}_i \theta_i + \bar{\theta}_i \eta_i)) \\ &= \int d\bar{\theta}_i d\theta_i (1 + i(\bar{\theta}, A\theta))(1 + i\bar{\eta}\theta)(1 + i\bar{\theta}\eta), \\ &= -i \det A \exp(-i\bar{\eta}_i A_{ij}^{-1} \eta_j). \end{aligned}$$

- Generalize to functional integrals over fermionic fields;

$$\begin{aligned} Z[\bar{\eta}, \eta] &= \int [d\bar{\psi}][d\psi] \exp(i \int d^4x [\bar{\psi}(i\cancel{\partial} - m)\psi + \bar{\eta}\psi + \bar{\psi}\eta]) \\ &= Z_0 \exp[- \int d^4x d^4y \bar{\eta}(x) S_F(x - y) \eta(y)]. \end{aligned}$$

where

$$S_F[x - y] = i(i\cancel{\partial} - m)^{-1} = \int \frac{d^4k}{(2\pi)^4} \frac{ie^{-ik(x-y)}}{\not{k} - m + i\epsilon}.$$

Get e.g.

$$\langle 0|T\psi(x)\bar{\psi}(y)|0\rangle = Z_0^{-1}(-i\frac{\delta}{\delta\bar{\eta}(x)})(i\frac{\delta}{\delta\eta(y)})Z[\eta, \bar{\eta}]|_{\eta, \bar{\eta}=0} = S_F(x - y).$$

Gives the Feynman rules for fermions that we discussed last quarter.

- For fermions, the $\det B$ is in the numerator, whereas for scalars it's in the denominator. The functional integral gives e^{iW} . So the sign of the contribution to W is opposite for closed scalar vs fermion loops: every closed fermion loop gets an extra -1 factor. (This relative minus sign is put to good use with supersymmetry!)

- Now gauge fields. Important point: gauge invariance.