

### 3/1/16 Lecture 17 outline

- Last time: Recall spin 1 gauge field canonical quantization, and  $A_\mu$  as an operator. Recall gauge invariance of  $\mathcal{L}_{EM}$ , needed to avoid having wrong sign kinetic term for longitudinal polarization terms.

- Functional integral for gauge fields. Important point: gauge invariance. Write  $A = A_\mu dx^\mu$ . Recall gauge symmetry  $A \rightarrow A^\alpha = A + d\alpha(x)$ , with  $\psi_i \rightarrow e^{-iq_i\alpha(x)}\psi$ . Redundancy in description, can only observe gauge invariant quantities. Need to replace  $\partial_\mu\psi_i \rightarrow D_\mu\psi_i \equiv (\partial_\mu + iq_i A_\mu)\psi_i$ . Then  $D_\mu^\alpha\psi_i^\alpha = e^{-iq_i\alpha}D_\mu\psi_i$  transforms nicely, with just an overall phase, and  $\bar{\psi}_i D_\mu\psi_i$  is gauge invariant. So the Dirac lagrangian,  $\bar{\psi}(i\not{D} - m)\psi$  is gauge invariant.

The terms linear in  $A_\mu$  give  $\mathcal{L} \supset -A_\mu j^\mu$ , with  $j^\mu$  the conserved current.

- In the functional integral, will have  $\int[dA] \exp(iS)$ . Integration measure must be gauge invariant, implies it gets a factor of gauge orbit volume. Would like to integrate only over a slice of inequivalent gauge fields, without integrating over the gauge orbits. Need to do this, since otherwise there is no well defined  $B^{-1}$ . Recall  $S = \int d^4x [-\frac{1}{4}F_{\mu\nu}^2] = \frac{1}{2} \int d^4k A_\mu(x)(\partial^2 g^{\mu\nu} - \partial^\mu\partial^\nu)A_\nu(x)$ . Note action vanishes if  $\tilde{A}_\mu(k) = k_\mu\alpha(k)$ . Gauge invariance.  $A_\mu^T = P_{\mu\nu}A^\nu$ ,  $P_{\mu\nu} = g_{\mu\nu} - \partial_\mu\partial_\nu/\partial^2$ .  $-\frac{1}{4}F_{\mu\nu}F^{\mu\nu} = \frac{1}{2}A_\mu^T\partial^2 g^{\mu\nu}A_\nu^T$ . Can't invert kinetic terms uniquely to find Green's function. We need to fix the gauge.

The functional integral should be over  $\int[dA^\mu]/(GE)$ , where we divide by the volume of the gauge equivalent orbits. The gauge equivalent orbits are associated with gauge transformations  $\alpha(x)$ , e.g.  $A_\mu \rightarrow A_\mu + \partial_\mu\alpha(x)$  in the Abelian case. We want to do the functional integral over  $A^\mu$ , dividing out by the  $\alpha(x)$ .

(Here are some details: Do this via

$$1 = \int [d\alpha(x)] \delta(G(A^\alpha)) \det\left(\frac{\delta G(A^\alpha)}{\delta\alpha}\right) = \Delta \int [d\alpha] \delta(G(A^\alpha)),$$

where  $G(A) = 0$  is some gauge fixing condition, e.g. Lorentz gauge,  $G(A) = \partial_\mu A^\mu$  and

$$\Delta = \det\left(\frac{\delta G(A^\alpha)}{\delta\alpha}\right)_{G=0}.$$

$\Delta$  is the Faddeev-Popov determinant. Write the functional integral as (using the gauge invariance of measure and action)

$$\int [d\alpha][dA] \Delta \delta(G[A]) \exp(iS[A]).$$

Have factored out the integral over the group volume. We can then just easily divide out by  $[d\alpha]$ , just cross it out. What's left is the gauge fixing delta function, and appropriate determinant factor.